



# $^2\text{H}$ and $^{139}\text{La}$ NMR Spectroscopy in Aqueous Solutions at Geochemical Pressures

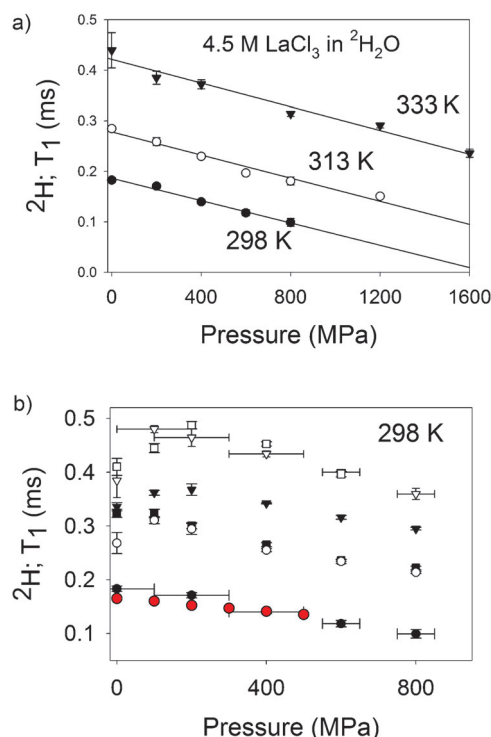
Gerardo Ochoa, Corey D. Pilgrim, Michele N. Martin, Christopher A. Colla, Peter Klavins, Matthew P. Augustine, and William H. Casey\*

**Abstract:** Nuclear spin relaxation rates of  $^2\text{H}$  and  $^{139}\text{La}$  in  $\text{LaCl}_3 + ^2\text{H}_2\text{O}$  and  $\text{La}(\text{ClO}_4)_3 + ^2\text{H}_2\text{O}$  solutions were determined as a function of pressure in order to demonstrate a new NMR probe designed for solution spectroscopy at geochemical pressures. The  $^2\text{H}$  longitudinal relaxation rates ( $T_1$ ) vary linearly to 1.6 GPa, consistent with previous work at lower pressures. The  $^{139}\text{La}$   $T_1$  values vary both with solution chemistry and pressure, but converge with pressure, suggesting that the combined effects of increased viscosity and enhanced rates of ligand exchange control relaxation. This simple NMR probe design allows experiments on aqueous solutions to pressures corresponding roughly to those at the base of the Earth's continental crust.

Geochemists invest much energy in models of aqueous solutions at elevated temperatures and pressures. A common model was recently extended to 1200 °C and 6.0 GPa<sup>[1]</sup> and predictions are now possible over such an extraordinary range of temperatures and pressures that experimental verification is difficult. This fact led us to design a simple NMR probe that allows experiments on solutions to 2.0 GPa.<sup>[2]</sup> Although there is a long history of high-pressure NMR in solution chemistry,<sup>[3]</sup> the pressures are usually limited to 0.5 GPa or less, with a few pioneering exceptions.<sup>[4]</sup>

Here we show that the design can reproduce previous results at lower pressures (< 0.5 GPa) and we extend them to nearly 2.0 GPa. We choose  $\text{LaCl}_3$  and  $\text{La}(\text{ClO}_4)_3$  electrolytes in  $^2\text{H}_2\text{O}$  because the Jonas group<sup>[5]</sup> previously measured  $^2\text{H}$   $T_1$  values on similar solutions, as well as viscosities, up to 0.5 GPa. Furthermore, these 1:3 electrolytes have a considerable freezing-point depression, which allows us to achieve gigapascal pressures at ambient temperatures.

We observe  $T_1$  values for  $^2\text{H}$  in  $\text{LaCl}_3 + ^2\text{H}_2\text{O}$  solutions to decrease nearly linearly with pressures up to 1.6 GPa (Figure 1). Results for 4.5 M  $\text{LaCl}_3$  solutions are shown in



**Figure 1.** a)  $T_1$  values from  $^2\text{H}$  NMR spectra as a function of pressure of 4.5 M  $\text{LaCl}_3 + ^2\text{H}_2\text{O}$  solutions. The lines are linear regressions. The  $T_1$  values as a function of solution composition at 298 K are shown in (b). The solid circles (●) are 4.5 M  $\text{LaCl}_3$  and the red symbols (●) identify data of Lee et al. (1974).<sup>[5]</sup> The  $\text{LaCl}_3 + ^2\text{H}_2\text{O}$  solutions are: ○ = 1.0 M, ▼ = 0.5 M, ▽ = 0.1 M, and the  $\text{La}(\text{ClO}_4)_3 + ^2\text{H}_2\text{O}$  solutions are: ■ = 1.0 M and □ = 0.1 M. Uncertainties in pressure are  $\pm 100$  MPa at 400 MPa or less and 50 MPa at pressures higher than 400 MPa. They are shown only for the 0.1 M and 4.5 M  $\text{LaCl}_3$  data to avoid clutter and are assigned as the 95 % prediction interval from repeated external calibrations (see the Supporting Information).

[\*] G. Ochoa, C. D. Pilgrim, M. N. Martin, Prof. M. P. Augustine, Prof. W. H. Casey  
Department of Chemistry  
University of California, Davis  
1 Shields Ave, Davis, CA 95616 (USA)  
E-mail: whcasey@ucdavis.edu

P. Klavins  
Department of Physics, University of California  
Davis (USA)

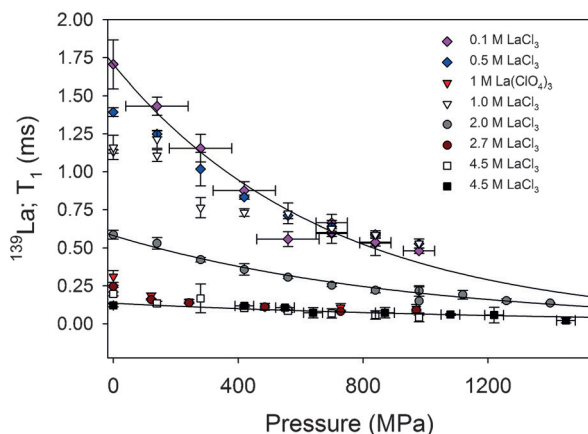
C. A. Colla, Prof. W. H. Casey  
Department of Earth and Planetary Sciences  
University of California, Davis (USA)

Supporting information for this article is available on the WWW under <http://dx.doi.org/10.1002/anie.201507773>.

Figure 1a and a full set of  $^2\text{H}$   $T_1$  measurements as a function of solution composition is shown in Figure 1b. The data overlay those of Lee et al.<sup>[5]</sup> at  $P$  smaller than 0.5 GPa and exhibit the same variation with pressure. In most experiments,  $^2\text{H}$   $T_1$  values initially rise, as was observed in previous work. The rise is attributed to electrolyte-induced solvent structuring<sup>[6,17]</sup> and values reach a maximum value at approximately 0.1–0.2 GPa.<sup>[3d,e,7]</sup> In the solutions with varying electrolyte concentrations, the  $T_1$  values decrease as pressures increase above 0.2 GPa (Figure 1b) undoubtedly reflecting increases in the solution viscosities, as was originally interpreted by Lee

et al.,<sup>[5]</sup> albeit over a smaller pressure range. Save possibly for the 4.5 M LaCl<sub>3</sub> solutions, only a fraction of the <sup>2</sup>H were held in solvation spheres of La<sup>III</sup>.

The  $T_1$  values for <sup>139</sup>La, however, differ considerably with concentration under ambient conditions and the values also vary much more with pressure than  $T_1$  values for <sup>2</sup>H in these same solutions (Figure 2). The differences in <sup>139</sup>La  $T_1$  values diminish with pressure until the values for all solutions



**Figure 2.** Rates of longitudinal relaxation ( $T_1$ ) for <sup>139</sup>La at 298 K in LaCl<sub>3</sub> + <sup>2</sup>H<sub>2</sub>O and La(ClO<sub>4</sub>)<sub>3</sub> + <sup>2</sup>H<sub>2</sub>O solutions. Uncertainties in  $T_1$  correspond to the estimated standard deviation of five replicate measurements. Uncertainties in pressure are usually  $\pm 100$  MPa at 400 MPa or less and 50 MPa at pressures higher than 400 MPa; they are shown only for the 0.1 M LaCl<sub>3</sub> data to avoid clutter. An exception is the 2.7 M LaCl<sub>3</sub> experiment where pressure was measured in situ and are assigned from the uncertainty in the spectral measurement ( $\pm 30$  MPa). See the Supporting Information for discussion of pressure uncertainties. Lines are exponential fits to three of the solution values.

converge to near 0.25 ms near  $P \gtrsim 1.2$  GPa, which also is close to the  $T_1$  values observed for the most concentrated solutions (Figure 2). As expected, acidification had little effect on the measured <sup>139</sup>La  $T_1$  values, indicating that the results are not strongly influenced by the hydrolysis of La<sup>III</sup> ( $pK_1 = 8.5^{[8]}$ ). The <sup>139</sup>La  $T_1$  value for an acidic solution of 1.0 M LaCl<sub>3</sub> at pD = 2.6, is  $1.02 \pm 0.03$  ms virtually identical to the unacidified sample (pD = 5.45;  $T_1 = 1.04 \pm 0.03$  ms).

There are two important terms affecting these relaxation rates. For most aqueous solutions in a high magnetic field, the time scale of tumbling of small molecules is shorter than the inverse Larmor frequency, which causes the longitudinal and transverse relaxation rates to be equal:  $1/T_1 = 1/T_2$ . In this “extreme-narrowing” limit the relaxation rate for quadrupolar nuclei in solution, such as <sup>139</sup>La ( $I = 5/2$ ) and <sup>2</sup>H ( $I = 1$ ) is given by Equation (1),<sup>[9]</sup>

$$\frac{1}{T_2} = \frac{1}{T_{10}} = \frac{3}{10} \pi^2 \frac{2I+3}{I^2(2I-1)} \langle \chi^2 \rangle \tau_c \quad (1)$$

where  $I$  is the spin of the nucleus in question,  $\tau_c$  is the correlation or reorientation time because of molecular

tumbling, and  $\langle \chi^2 \rangle$  is the mean square of the zero-average, fluctuating quadrupolar coupling constant. Values of the quadrupolar coupling constants are known for <sup>2</sup>H and <sup>139</sup>La nuclei. In LaCl<sub>3</sub> + <sup>2</sup>H<sub>2</sub>O solution at 20 °C,  $\langle \chi \rangle_{rms} = 237.7 \pm 10$  kHz<sup>[10]</sup> and  $\langle \chi \rangle_{rms} = 3.1\text{--}3.4$  MHz<sup>[10e,11]</sup> values for <sup>2</sup>H and <sup>139</sup>La have been calculated from experimental data and Equation (1). Lang, Ludemann, and co-workers<sup>[3f]</sup> report  $\langle \chi \rangle_{rms} = 230$  MHz for <sup>2</sup>H up to 0.9 GPa and noted that it increases to  $\langle \chi \rangle_{rms} = 264$  MHz in 11 M LiCl solutions. The quadrupolar coupling constant  $\langle \chi \rangle_{rms}$  for <sup>139</sup>La in La(ClO<sub>4</sub>)<sub>3</sub> solutions is identical to that in LaCl<sub>3</sub> solutions at ambient pressures.<sup>[11]</sup>

The correlation time,  $\tau_c$ , is strongly affected by solution viscosity, which modulates the rates of random diffusive movements and varies considerably with pressure. The link between viscosities and correlation times is usually treated using a version of the Stokes–Einstein Equation (2),<sup>[5]</sup>

$$\tau_c = \frac{\eta \kappa V_m}{kT} \quad (2)$$

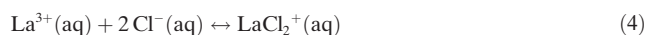
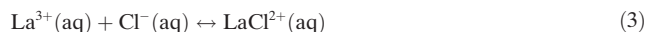
where  $V_m$  is the hydrodynamic volume and  $\eta$  is the macroscopic viscosity. This approach ignores details of the nuclear spin relaxation pathways but is appropriate for the experiments described here. The parameter  $\kappa$  describes coupling of translation and rotation during relaxation<sup>[13]</sup> and is usually set to unity so that  $V_m$  can be treated as the effective hydrodynamic volume for the <sup>2</sup>H<sub>2</sub>O molecule, or the solvated La<sup>III</sup> complex. Since viscosity increases with pressure, the time scale of molecular tumbling changes with pressure and causes the value of  $\tau_c$  to increase.

For most solutions in this study, the relaxation time of <sup>2</sup>H reflects only nuclei in bulk solution and not solvation spheres. For solutions containing structure-making cations such as La<sup>3+</sup>(aq), the rotation of <sup>2</sup>H<sub>2</sub>O molecules that are tightly bound in solvation spheres has only a small effect on the local structure of the liquid.<sup>[14]</sup> Chloride ions have been shown to have little effect on the effective volume of the <sup>2</sup>H<sub>2</sub>O molecule.<sup>[15]</sup> Changes in viscosity with pressure exert the dominant controls on relaxation. The  $T_1$  values for <sup>2</sup>H in various LaCl<sub>3</sub> solutions vary by roughly a factor of 2–3 as a function of electrolyte concentration (Figure 1b) and this range can bind the effects of viscosity changes alone over our pressure range.

However, for <sup>139</sup>La relaxation,  $T_1$  relaxation rates are affected by viscosity but also from quadrupolar interactions of a small set of La<sup>III</sup> complexes: La<sup>3+</sup>(aq), LaCl<sup>2+</sup>(aq), and LaCl<sub>2</sub><sup>+</sup>(aq). The <sup>139</sup>La  $T_1$  values, in contrast to those of <sup>2</sup>H, vary more with solution composition and pressure (Figure 2), suggesting that viscosity changes alone do not fully account for the variation in <sup>139</sup>La  $T_1$  values. The effects of concentration on <sup>139</sup>La relaxation clearly diminish with increased pressure. Furthermore, if the correlation times of the two nuclei, <sup>139</sup>La and <sup>2</sup>H, are controlled only by a single bulk variable, such as solution viscosity, the ratio of  $T_1$  values for <sup>2</sup>H and <sup>139</sup>La should be constant and independent of pressure. A plot of the ratio of these values for the 2.7 M LaCl<sub>3</sub> solution (see Figure S4 in the Supporting Information) is constant to within 10% at  $P > 0.6$  GPa indicating that the <sup>139</sup>La NMR relaxation is well correlated with <sup>2</sup>H relaxation at higher

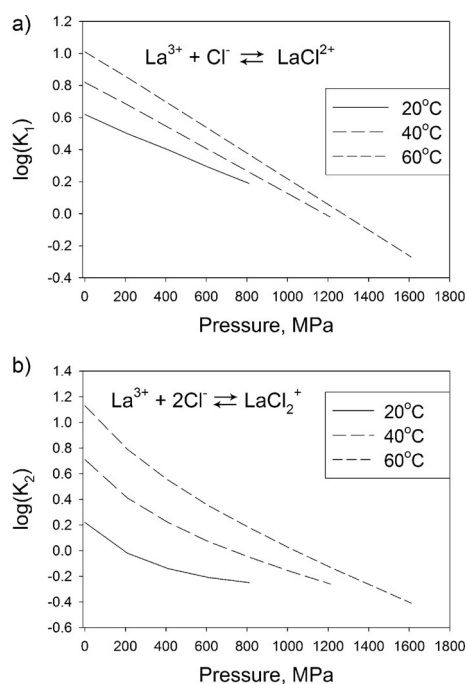
pressures. Clearly viscosity is a dominating variable and becomes increasingly important at higher pressures.

The trend toward convergence of  $^{139}\text{La}$   $T_1$  values beyond  $P \gtrsim 1.2$  GPa indicates that the product ( $\chi^2 \cdot \tau_c$ ) becomes progressively insensitive to solution composition at elevated pressures. Any process that changes the electric field gradient at the  $^{139}\text{La}$  nucleus, and these include ion pairing or rapid ligand exchanges, could affect relaxation through the quadrupolar coupling constant. At the time when the Lee et al. paper<sup>[5]</sup> was published (1974), there was no firm evidence for complexation of  $\text{La}^{3+}$  and  $\text{Cl}^-$  although the authors clearly knew of this possibility. There have since been careful NMR,<sup>[10e,16]</sup> X-ray,<sup>[17]</sup> vibrational,<sup>[18]</sup> and other spectroscopic<sup>[19]</sup> studies of the reactions given in Equations (3) and (4)



and these are supported by theoretical studies.<sup>[20]</sup> Equilibrium constants ( $K_1$ ;  $K_2$ ) for reactions (3) and (4) have been tabulated<sup>[21]</sup> and Migisidov, Williams, and Wagner<sup>[22]</sup> organized the observations into parameters for the Helgeson–Kirkham–Flowers (HKF) model, allowing predictions to elevated temperatures and pressures.

Values of  $K_1$  and  $K_2$  are shown as a function of pressure in Figure 3. The equilibrium constants are small, consistent with



**Figure 3.** Estimated equilibrium constants for a) reaction (3) and b) reaction (4).<sup>[22]</sup>

weak outer-sphere ion-pairing reaction, and diminish further with increased pressure. The trend is consistent with the favoring of more highly solvated ions with pressure because solvent water molecules can be packed into a smaller volume.

There are reports of ion pairing for  $\text{La}(\text{ClO}_4)_3$  solutions<sup>[18]</sup> but these constants have not been parameterized for elevated pressures. Ion pairing would be weaker and yet would undoubtedly trend the same as the chloride data.

However, a large fraction of the  $\text{La}^{\text{III}}$  is present as  $\text{LaCl}^{2+}$  and  $\text{LaCl}_2^+$  at all pressures. A quick calculation illustrates the magnitude—the equilibrium constants for Equations (3) and (4) are  $K_1 = 10^{0.19}$  and  $K_2 = 10^{-0.25}$ , at pressures of 0.80 GPa. At this pressure, 77% of  $\text{La}^{\text{III}}$  is present as  $\text{La}^{3+}(\text{aq})$  with the remainder as  $\text{LaCl}^{2+}(\text{aq})$  and  $\text{LaCl}_2^+(\text{aq})$ . The estimates are crude but show ion pairing even at the highest pressures. (For the calculations we use the extended Debye–Hückel equation for activity corrections in our least concentrated solution, 0.1 M  $\text{LaCl}_3$ , with  $A = 0.5042$ ,  $a = 5.0$  Å for  $\text{La}^{\text{III}}$  ions,  $a = 3.0$  Å for  $\text{Cl}^-$ ,  $B = 0.3288$  ( $\text{kg mol}^{-1}$ ) $^{-0.5}$  Å and  $B = 0.0410$  kmol $^{-1}$ . Choice of another model for activity corrections makes no essential difference.) The key result is that interconversion of  $\text{La}^{3+}$ ,  $\text{LaCl}^{2+}$ , and  $\text{LaCl}_2^+$  is not negligible even at high pressures.

Our observation that  $^{139}\text{La}$   $T_1$  values converge at  $P \gtrsim 1.2$  GPa indicates that the electric field gradients at the  $^{139}\text{La}$  nuclei are becoming uniform even in the presence of large speciation differences. Viscosity effects are clearly dominating, but such a case would arise if all  $\text{La}^{\text{III}}$  species interconvert by rapid rates of ligand substitution at very high pressures. If such rapid rates were occurring, the exchanging ligands must be dominated by solvation waters because we observe limited influence of changes in anion concentrations. Rates of solvent exchange for  $\text{La}^{3+}$  fall into the range of  $10^{8.5} - 10^9$  s $^{-1}$ ,<sup>[23]</sup> which is near the Smoluchowski limit of diffusion-controlled reactions and cannot reach the  $10^{-12}$  s timescales of nuclear reorientation. More likely is that rapid ligand substitutions diminish the electric field gradients, and thus  $\chi$ , at high pressures.

The major goal of this work was to show that a simple new NMR probe design can be used for solute spectroscopy at geochemical pressures. Uncertainties in pressure can also be reduced to a few tens of megapascals when the ruby is included in the cell, as is illustrated in the table-of-contents picture, because this geometry reduces the error in pressure to the precision of the fluorescence wavelength measurement, rather than the statistics associated with use of an external calibration scale.

## Acknowledgements

This work is supported by the Department of Energy (grant number DE-FG02-05ER15693). Additional funding sources are listed in the Supporting Information. The authors thank Drs. Jeff Walton and Brent Pautler for assistance, Profs. Sverjensky and Manning for encouragement and Mr. Brian Devine for design suggestions.

**Keywords:** aqueous solutions · geochemistry · high pressure · NMR probe design · NMR spectroscopy

**How to cite:** *Angew. Chem. Int. Ed.* **2015**, *54*, 15444–15447  
*Angew. Chem.* **2015**, *127*, 15664–15667

- [1] a) D. A. Sverjensky, B. Harrison, D. Azzolini, *Geochim. Cosmochim. Acta* **2014**, *129*, 125–145; b) D. Pan, L. Spanu, B. Harrison, D. A. Sverjensky, G. Galli, *Proc. Natl. Acad. Sci. USA* **2013**, *110*, 6646–6650.
- [2] B. G. Pautler, C. A. Colla, R. L. Johnson, P. Klavins, S. J. Harley, C. A. Ohlin, D. A. Sverjensky, J. W. Walton, W. H. Casey, *Angew. Chem. Int. Ed.* **2014**, *53*, 9788–9791; *Angew. Chem.* **2014**, *126*, 9946–9949.
- [3] a) A. Drljaca, C. D. Hubbard, R. Van Eldik, T. Asano, M. V. Basilevsky, W. J. Le Noble, *Chem. Rev.* **1998**, *98*, 2167–2289; b) R. Van Eldik, *Advances in Inorganic Chemistry*, Vol. 62, Academic Press, San Diego, **2010**; c) R. Van Eldik, T. Asano, W. J. Le Noble, *Chem. Rev.* **1989**, *89*, 549–688; d) E. W. Lang, H. D. Ludemann, *J. Chem. Phys.* **1977**, *67*, 718–723; e) E. W. Lang, D. Girlich, H. D. Ludemann, L. Piculell, D. Muller, *J. Chem. Phys.* **1990**, *93*, 4796–4803; f) E. W. Lang, H. D. Ludemann, *Prog. Nucl. Magn. Reson. Spectrosc.* **1993**, *25*, 507–633.
- [4] a) L. Ballard, C. Reiner, J. Jonas, *J. Magn. Reson. Ser. A* **1996**, *123*, 81–86; b) M. de Langen, K. O. Prins, *Rev. Sci. Instrum.* **1995**, *66*, 5218–5222.
- [5] Y. K. Lee, J. H. Campbell, J. Jonas, *J. Chem. Phys.* **1974**, *60*, 3537–3543.
- [6] Y. Ducommun, W. L. Earl, A. E. Merbach, *Inorg. Chem.* **1979**, *18*, 2754–2758.
- [7] a) E. W. Lang, H.-D. Ludemann, *Angew. Chem. Int. Ed. Eng.* **1982**, *21*, 315–329; b) Y. Lee, J. Jonas, *J. Magn. Reson. (1969-1992)* **1971**, *5*, 267–272.
- [8] C. F. Baes, R. E. Mesmer, *The Hydrolysis of Cations*, Wiley, New York, **1976**.
- [9] R. K. Harris, *Nuclear Magnetic Resonance Spectroscopy-A Physicochemical View*, Longman Scientific and Technical, Essex, U.K., **1986**.
- [10] a) H. W. Spiess, B. B. Garrett, R. K. Sheline, S. W. Rabideau, *J. Chem. Phys.* **1969**, *51*, 1201–1205; b) D. E. Woessner, *J. Chem. Phys.* **1964**, *40*, 2341–2348; c) J. Powles, M. Rhodes, J. Strange, *Mol. Phys.* **1966**, *11*, 515–523; d) B. B. Garrett, A. B. Denison, S. W. Rabideau, *J. Phys. Chem.* **1967**, *71*, 2606–2611; e) T. Yaita, D. Ito, S. Tachimori, *J. Phys. Chem. B* **1998**, *102*, 3886–3891.
- [11] J. Reuben, *J. Phys. Chem.* **1975**, *79*, 2154–2157.
- [12] D. Laage, G. Stirnemann, F. Sterpone, R. Rey, J. T. James, T. Hynes, *Annu. Rev. Phys. Chem.* **2011**, *62*, 395–416.
- [13] a) D. Kivelson, M. G. Kivelson, I. Oppenheim, *J. Chem. Phys.* **1970**, *52*, 1810–1821; b) R. E. D. McClung, D. Kivelson, *J. Chem. Phys.* **1968**, *49*, 3380–3391.
- [14] O. V. Yazyev, L. Helm, *J. Chem. Phys.* **2006**, *125*, 054503.
- [15] H. G. Hertz, in *Water, A Comprehensive Treatise; Aqueous Solutions of Simple Electrolytes*, Vol. 3, Springer, Plenum, New York, **1973**.
- [16] Z. Chen, C. Detellier, *J. Solution Chem.* **1992**, *21*, 941–952.
- [17] S. Díaz-Moreno, S. Ramos, D. T. Bowron, *J. Phys. Chem. A* **2011**, *115*, 6575–6581.
- [18] W. W. Rudolph, G. Irmer, *Dalton Trans.* **2015**, *44*, 295–305.
- [19] V. Sharma, F. Boehm, M. Seitz, G. Schwaab, M. Havenith, *Phys. Chem. Chem. Phys.* **2013**, *15*, 8383–8391.
- [20] a) M. Bühl, N. Sieffert, A. Partouche, A. Chaumont, G. Wipff, *Inorg. Chem.* **2012**, *51*, 13396–13407; b) P. X. Rutkowski, M. C. Michelini, J. K. Gibson, *Phys. Chem. Chem. Phys.* **2012**, *14*, 1965–1977.
- [21] a) S. A. Wood, *Chem. Geol.* **1990**, *82*, 159–156; b) F. J. Millero, *Geochim. Cosmochim. Acta* **1992**, *56*, 3123–3133.
- [22] A. A. Mignidov, A. E. Williams-Jones, T. Wagner, *Geochim. Cosmochim. Acta* **2009**, *73*, 7087–7109.
- [23] L. Helm, A. E. Merbach, *Chem. Rev.* **2005**, *105*, 1923–1959.

Received: August 19, 2015

Published online: September 25, 2015